# A WIMPy Baryogenesis Miracle

Baryogenesis via WIMP annihilation

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University of Chicago Theory Seminar

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#### Outline

- Motivation: WIMP miracle and dark matter/baryon ratio
- Review of baryogenesis
- Example: WIMPy leptogenesis
- WIMP annihilation to quarks
- Constraints and detection prospects

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    - \* The WIMP miracle
  - ② Dark matter/baryon ratio:  $\Omega_{\rm DM} \approx 5 \, \Omega_{\rm baryon}$
- Our models incorporate both observations
  - Oark matter abundance: Established by thermal freeze-out according to the WIMP miracle
  - Oark matter/baryon ratio: Dark matter annihilation generates a baryon asymmetry
    - ★ Connection between the dark and visible sector abundances
- For a model incorporating the WIMP miracle in baryogenesis in a different way than WIMPy baryogenesis, see McDonald, 1009.3227 and 1108.4653

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 Relic abundance inversely proportional to annihilation cross section

$$\Omega_{\rm WIMP} \approx \Omega_{\rm DM} \, \frac{1 \, \, {\rm pb}}{\langle \sigma_{\rm ann} \, v \rangle}$$

- In WIMP miracle framework,  $\Omega_{\rm DM} \sim \Omega_{\rm baryon}$  is a coincidence
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- Can we have some features of symmetric dark matter while also establishing a connection between the dark matter and baryon abundances?

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- WIMPy baryogenesis is nice because it
  - ► Ties all dark matter and baryogenesis physics to the weak scale
    - \* Possible weak scale origin of new fields and couplings?
  - ► Gives indirect detection signals of conventional symmetric WIMP dark matter
  - Incorporates baryogenesis by annihilation, which has often been overlooked
    - \* Proposed by Bento, Berezhiani 2001; Gu, Sarkar 2009

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  - Violation of baryon number
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  - 3 Departure from thermal equilibrium
- All three conditions are satisfied in the Standard Model but
  - ▶ CP violation not big enough (suppressed by 12 Yukawa couplings  $\sim 10^{-20}$ )
  - Phase transition not first order

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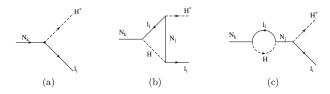
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- lacksquare B or L violation: Majorana mass of RH neutrino violates L, lepton asymmetry transferred to B by sphalerons
- $\bigcirc$  *CP* violation: *CP*-violating phases in  $u_{\nu}$
- $\odot$  Departure from equilibrium: N decays out of equilibrium
- If we only considered tree level diagram, CP phases disappear with  $|\mathcal{M}|^2$ 
  - Need to consider interference of tree and loop diagrams



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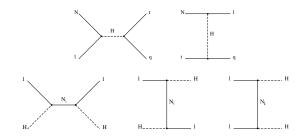
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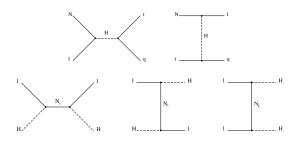
$$\sim \frac{1}{4\pi} \frac{\text{Im}(y_{\nu \, 11}^* y_{\nu \, i1}^* y_{\nu \, ij} y_{\nu \, 1j})}{|y_{\nu \, 11}|^2} \frac{m_{N1}}{m_{Ni}}$$

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- Asymmetry proportional to number of decays that happen after washout freezes out (at  $T \ll m_{N1}$ )
  - ▶  $N_1$  lifetime longer than Hubble time at  $T = m_{N1}$  ( $\Gamma_{N1} < H(m_{N1})$ )

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- How can washout go out of equilibrium sufficiently early?
  - One of lepton-number-carrying fields is heavy or washout cross section much smaller than annihilation cross section

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- Lepton asymmetry transferred to baryon asymmetry by sphalerons
  - ▶ Sphalerons ineffective after electroweak phase transition  $(T_c \sim 100 \; \mathrm{GeV})$
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- ullet New field  $\psi$ 
  - $\psi$  is a doublet with hypercharge +1/2
  - ightharpoonup To allow the widest possible range of masses, take  $\psi$  to be vectorlike

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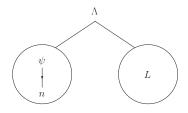
- $\bullet~{\rm U}(1)$  symmetry under which  $L,~\psi$  oppositely charged
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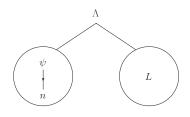
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- Two possible solutions:
  - Two sectors with separately preserved asymmetries
    - **\*** Simplest  $\psi$  decay:  $\psi \to H n$ , where n is a singlet
  - $\bullet$   $\psi$  decays with U(1)-violating couplings



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- Minimal solution:  $Z_4$  symmetry
  - ▶ Charge of X = i
  - $\qquad \qquad \textbf{ Charge of } \psi = -1$
  - ▶ Charge of SM fields = +1
- ullet Since X has a  $Z_4$  charge, it must be Dirac

#### A minimal "complete" model:

• We choose the simplest UV completion: effective operator arises from exchange of pseudoscalars  $S_{\alpha}$ 

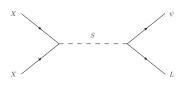
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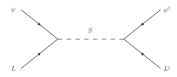
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Annihilation and washout scatterings:



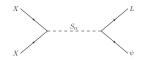
$$\sigma_{\rm ann} \sim |\lambda_X|^2 |\lambda_L|^2$$

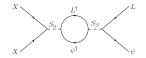


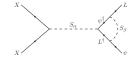
$$\sigma_{\rm washout} \sim |\lambda_L|^4$$

lacktriangle Baryon number violation  $\checkmark$ 

- Baryon number violation
- OP violation:
  - ▶ CP phases in couplings  $\lambda_X$ ,  $\lambda_L$
  - Interference of tree and loop diagrams





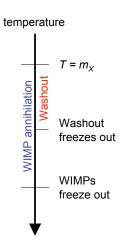


 $\bullet$  Must have at least two generations of S for non-zero CP phase in amplitude

- Baryon number violation
- $\bigcirc$  CP violation  $\checkmark$

- Baryon number violation √
- ② CP violation ✓
- Departure from thermal equilibrium?

- Asymmetry generated while DM annihilates
- Washout eliminates asymmetry as it accumulates
- Need to have washout freeze out during era of rapid WIMP annihilation



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- At  $z \gg 1$ , X and lepton asymmetry ( $\Delta L$ ) go out of equilibrium
  - ▶ Determines WIMP relic abundance and baryon asymmetry

• Evolution of  $Y_X$  and  $Y_{\Delta L}$ :

#### **Boltzmann equations:**

$$\frac{dY_a}{dz} = -\frac{(2\pi)^4}{z H(z) s(z)} \int d\Pi_a d\Pi_b d\Pi_c d\Pi_d |\mathcal{M}_{ab\to cd}|^2 \delta^4(\sum p) (f_a f_b - f_c f_d)$$

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  - ▶ Collision term proportional to annihilation cross section  $\langle \sigma_{XX \to L\psi} v \rangle$
  - lacktriangle Drives  $Y_X$  to equilibrium value when scattering is rapid
  - ▶ Proportional to the *square* of the *X* distribution

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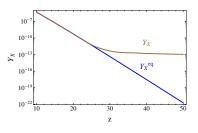
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$$\frac{dY_X}{dz} \sim -\langle \sigma_{XX \to L\psi} \, v \rangle \left[ Y_X^2 - (Y_X^{\text{eq}})^2 \right]$$

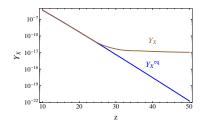
ullet If the baryon asymmetry is small, there is no back-reaction on  $Y_X$ 

• We get the conventional WIMP equation

• 
$$Y_X(z=\infty) \sim 1/\langle \sigma_{XX\to L\psi} v \rangle$$



- We get the conventional WIMP equation
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- For z>1, we want  $dY_X/dz\approx dY_X^{\rm eq}/dz$  if X tracks its equilibrium distribution
  - ▶ This implies a departure of X from thermal equilibrium!
  - Integrating the deviation from equilibrium over z gives  $\Delta Y_X$ , the total number of DM particles annihilated

### WIMPy leptogenesis: lepton asymmetry evolution

- Lepton asymmetry evolution:
  - ► Two important terms:
    - \* Asymmetry generation by XX annihilation (proportional to fractional asymmetry per annihilation  $\epsilon$ )
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$$\frac{dY_{\Delta L}}{dz} \sim +\epsilon \times (\text{WIMP ann. rate}) - Y_{\Delta L} \times (\text{washout rate})$$

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$$\sim -\epsilon \frac{dY_X}{dz} - Y_{\Delta L} \langle \sigma_{L\psi \to L^\dagger \psi^\dagger} \, v \rangle Y_L^{\text{eq}} \, Y_\psi^{\text{eq}}$$

$$\frac{dY_{\Delta L}}{dz} \sim -\epsilon \frac{dY_X}{dz} - Y_{\Delta L} \langle \sigma_{L\psi \to L^{\dagger}\psi^{\dagger}} v \rangle Y_L^{\text{eq}} Y_{\psi}^{\text{eq}}$$

 While annihilation is occurring, there is competition between asymmetry generation and washout

$$\frac{dY_{\Delta L}}{dz} \sim -\epsilon \, \frac{dY_X}{dz} - Y_{\Delta L} \langle \sigma_{L\psi \to L^\dagger \psi^\dagger} \, v \rangle \, Y_L^{\rm eq} \, Y_\psi^{\rm eq} \label{eq:equation_for_property}$$

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  - ► Early times: there is an instantaneous steady-state solution found by balancing the rates of asymmetry creation and depletion

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- ★ Too small for observed baryon asymmetry
- **Late times:** define  $z_0$  as the time when washout processes freeze out
  - ★ We're left with the equation

$$\frac{dY_{\Delta L}}{dz} \sim -\epsilon \frac{dY_X}{dz} \qquad (z > z_0)$$

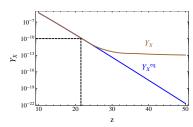
$$Y_{\Delta L}(\infty) \approx \epsilon \left[ Y_X(z_0) - Y_X(\infty) \right]$$

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- Washout must freeze out before annihilations cease
- ullet Washout freezes out when washout rate  $\lesssim$  Hubble scale
- Washout rate  $\sim \langle \sigma_{L\psi \to L^\dagger \psi^\dagger} \, v \rangle \, Y_L^{\mathrm{eq}} \, Y_\psi^{\mathrm{eq}}$

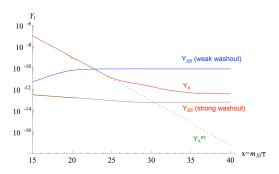
$$Y_{\Delta L} \sim 10^{-10}$$
 and  $\epsilon < 1 \implies z_0 \lesssim 20$ 

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  - **1** Heavy  $m_{\psi}$  so that  $Y_{\psi}^{\mathrm{eq}}$  is exponentially suppressed

# WIMPy leptogenesis: asymmetry

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- ullet Washout freezes out before WIMPs o weak washout
- ullet Washout freezes out after WIMPs o strong washout

# WIMPy leptogenesis

#### Recap so far:

- Baryogenesis through WIMP annihilation is possible if
  - ► Annihilation occurs through *L*-violating coupling
  - ▶ Non-zero *CP* phases in *L*-violating coupling
- Need washout to freeze out while WIMP annihilation is still active
- WIMPs described by equilibrium distribution during this time!

• CP-violating factor: fractional asymmetry generated by each annihilation

$$\epsilon = \frac{\sigma(XX \to \psi_i L_i) - \sigma(XX \to \psi_i^{\dagger} L_i^{\dagger})}{\sigma(XX \to \psi_i L_i) + \sigma(XX \to \psi_i^{\dagger} L_i^{\dagger})}$$

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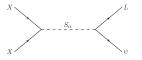
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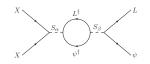
• Many free parameters! Make assumptions to include minimal ingredients, simplify analysis:

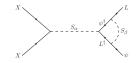
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- Many free parameters! Make assumptions to include minimal ingredients, simplify analysis:
  - ▶ Only one flavour of lepton relevant for WIMPy leptogenesis
  - lacktriangle Annihilation through the lightest scalar  $S_1$  is dominant
  - Phases are large

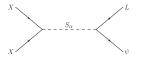


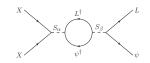


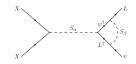


• With these assumptions:

$$\epsilon \sim -\frac{1}{4\pi} \, \frac{\mathrm{Im}(\lambda_{L1}^2 \lambda_{L2}^{*2})}{|\lambda_{L1}^2|} \, \frac{(2m_X)^2}{m_{\mathrm{S2}}^2} f\left(\frac{m_\psi}{2m_X}\right)$$





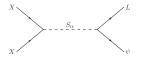


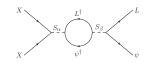
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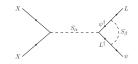
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• The requirement of dominant scattering through  $S_1$  (assume  $\sigma_{S2} < 0.2\sigma_{S1}$ ) gives a bound on  $\epsilon$ :

$$|\epsilon| \lesssim \frac{2\lambda_{L1}^2 m_X^2}{3\pi \sqrt{5} m_{S1}^2} \, f\left(\frac{m_\psi}{2m_X}\right)$$







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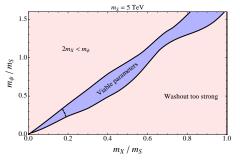
- $\bullet$  Masses and couplings of heavy  $S_\alpha$  contribute only indirectly through loop effects to  $\epsilon$ 
  - Use  $\epsilon$  as a free parameter, subject to bound

### Numerical results: masses

- ullet 6 parameters:  $m_X$ ,  $m_\psi$ ,  $m_S$ ,  $\lambda_X$ ,  $\lambda_L$ , and  $\epsilon$
- Show masses for which WIMPy leptogenesis gives correct relic density and asymmetry with perturbative couplings  $\lambda_L$ ,  $\lambda_X$ , and  $\epsilon$

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• X and  $\psi$  mass constrained to lie close together (within  $m_{\psi} \sim 1 - 2 m_{X}$ )

- $m_S = 5 \text{ TeV}$
- ullet Asymmetry should be generated before sphalerons decouple  $\Rightarrow m_X \gtrsim {\sf TeV}$ 
  - ▶ Dashed line in figure for Standard Model electroweak phase transition

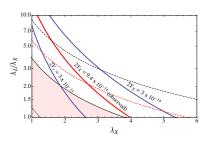
# Numerical results: couplings

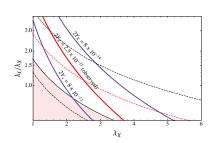
- Choose points in middle of parameter space:
  - $m_S=5$  TeV for both plots

# Numerical results: couplings

- Choose points in middle of parameter space:
  - $m_S = 5$  TeV for both plots
  - $m_X=4.25$  TeV,  $m_\psi=7.5$  TeV, and  $\epsilon=0.075$

•  $m_X=1.5$  TeV,  $m_\psi=2.25$  TeV, and  $\epsilon=0.0075$ 





- Solid lines: X relic abundance
- Dotted lines: baryon asymmetry (from top,  $Y_{\Delta B}=3\times 10^{-11}$ ,  $8.85\times 10^{-11}$ ,  $3\times 10^{-10}$ )
- Shaded region inconsistent with assumptions

- Constructed a concrete model of leptogenesis through WIMP annihilation
- Get correct WIMP relic density and baryon asymmetry with:
  - ▶ All masses  $\mathcal{O}(\text{TeV})$
  - All couplings  $\gtrsim 1$
  - lacktriangle Sufficiently large asymmetry in region with  $m_X \sim m_\psi$
- Limitation:  $T_{\text{lepto}} > T_{\text{electroweak}}$

### Annihilation to quarks

- Consider model similar to leptogenesis
  - lackbox WIMP annihilation to up quark ar u;  $\psi$  is colour triplet with charge +2/3

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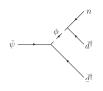
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- $\psi$  decays through operator  $\bar{\psi} \bar{d} \bar{d} n / \Lambda^2$  to quarks, singlet n
  - lacktriangle ex. decay through coloured scalar  $\phi$

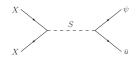
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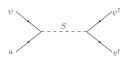


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$$\Delta \mathcal{L} = +\frac{i}{2} \left( \lambda_{X\alpha} X^2 + \lambda'_{X\alpha} \bar{X}^2 \right) S_{\alpha} + i \lambda_{B\alpha i} S_{\alpha} \bar{u}_i \psi_i + \text{h.c.}$$



$$\sigma_{\rm ann} \sim |\lambda_X|^2 |\lambda_B|^2$$



$$\sigma_{\rm washout} \sim |\lambda_B|^4$$

### Annihilation to quarks: numerical results

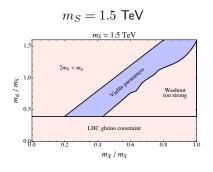
 $\bullet$  6 parameters:  $m_X$  ,  $m_\psi$  ,  $m_S$  ,  $\lambda_X$  ,  $\lambda_B$  , and  $\epsilon$ 

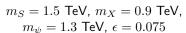
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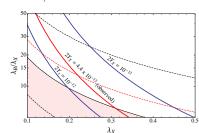
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  - $m_{\psi} \gtrsim 590 \text{ GeV}$
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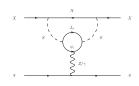
### Constraints and signals

- We consider (briefly) the three most important constraints/observable effects:
- Direct detection
- Indirect detection
- Colliders

# Constraints and signals: direct detection

#### Annihilation to leptons:

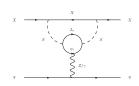
- Only couples to nucleons through 2-loop moment
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# Constraints and signals: direct detection

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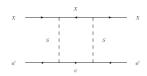
- Only couples to nucleons through 2-loop moment
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#### Annihilation to quarks:

Couples at one-loop:

$$\sigma_{X-N} \sim \frac{1}{16\pi} \left(\frac{\lambda_B^2 \lambda_X^2}{16\pi^2}\right)^2 \, \frac{\mu^2}{m_X^4} \label{eq:sigmaX}$$



- Benchmark points:
  - ①  $m_X=4.25$  TeV,  $m_\psi=7.25$  TeV,  $m_S=5$  TeV,  $\lambda_X=2.7$  and  $\lambda_B=4.5$ :  $\sigma_{X-N}\approx 1\times 10^{-44}~{\rm cm}^2$
  - ②  $m_X=0.9$  TeV,  $m_\psi=1.2$  TeV,  $m_S=1.5$  TeV,  $\lambda_X=0.22$  and  $\lambda_B=2.8$ :  $\sigma_{X-N}\approx 4\times 10^{-46}~{\rm cm}^2$

### Constraints and signals: indirect detection

- Both scenarios annihilate to quarks
- Best prospect for indirect detection: antideuterons
  - Very low astrophysical backgrounds at low energies
  - Donato, Fornengo, Salati 2000; Baer, Profumo 2005; Cui, Mason, Randall 2010

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### Annihilation to leptons:

- $XX \to W^{\pm}$ , h
- Hadronization in boosted frame
- Mass constraint reach  $\mathcal{O}(100~{\rm GeV})$

### Annihilation to quarks:

- $XX o \text{color-connected } \bar{u}\bar{d}\bar{d}$
- Some hadronization in rest frame
- Low-energy antideuterons!
- ullet Can exclude up to  $m_X \sim {\sf TeV}$

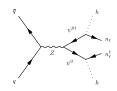
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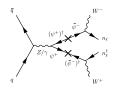
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### Leptogenesis:

$$\mathcal{L} \supset \lambda_i' \, \psi \, n \, H^{\dagger}$$



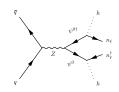


- ullet Strongest bound comes from chargino searches at LEP  $(m_\psi \gtrsim 100 \ {
  m GeV})$ 
  - $\tilde{\chi}^{\pm} \to W^{\pm} \tilde{\chi}^0 \to jj \, \tilde{\chi}^0$

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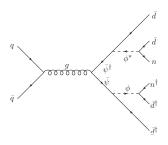


$$\bar{q}$$
 $W^ \bar{\psi}^ Z/\gamma_{\psi^+}$ 
 $(\bar{\psi}^-)^{\dagger}$ 
 $N_{\ell}$ 
 $W^+$ 

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  - $\tilde{\chi}^{\pm} \to W^{\pm} \, \tilde{\chi}^0 \to jj \, \tilde{\chi}^0$
- LHC not yet sensitive to electroweak production
  - lacktriangle May be able to find in targeted searches: b-tagging, reconstruct Higgs mass

Annihilation to quarks:

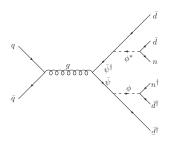
#### Annihilation to quarks:



$$\mathcal{L} \supset \lambda_i \, \bar{\psi}_i \, \bar{d}_i \, \phi^* + \lambda_i' \, \phi \, \bar{d}_i \, n_i$$

- Gluino-like topology with different group theory factors
- $4j + \cancel{E}_{\mathrm{T}}$  final state
- Current LHC bound excludes  $m_{\psi} \lesssim 590 \text{ GeV}$

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ullet LHC should test  $m_{\psi}$  up to  $\sim 2$  TeV at  $100~{
m fb}^{-1}$ 

#### Conclusions

- WIMPy baryogenesis: WIMP annihilations can generate a baryon asymmetry
- $\bullet$  Can get correct relic density and baryon asymmetry with  $\sim$  TeV masses,  $\mathcal{O}(1)$  couplings
  - ▶ Need  $m_X \sim m_\psi$
- Baryon asymmetry generated at weak scale (directly or via leptogenesis)
- Examined possible signals at the LHC and in dark matter detection experiments

# Back-up slides

Back-up slides

### WIMPy leptogenesis: Boltzmann equations

Evolution of the asymmetry in one **component** of the L doublet:

$$\begin{split} \frac{H(m_X)}{z} \, \frac{dY_X}{dz} &= & -4s (\sigma_{XX \to L_i \psi_i} \, v) [Y_X^2 - (Y_X^{\text{eq}})^2] - 2s \epsilon_X \frac{\xi \, Y_{\Delta L_i}}{Y_{\gamma}} (\sigma_{XX \to L_i \psi_i} \, v) (Y_X^{\text{eq}})^2 \\ &- \text{Br}_X^2 \langle \Gamma_S \rangle Y_S^{\text{eq}} \left( \frac{Y_X}{Y_X^{\text{eq}}} \right)^2 + \text{Br}_X \langle \Gamma_S \rangle \left( Y_S - \text{Br}_L \, Y_S^{\text{eq}} \right) - \epsilon \frac{\xi \, Y_{\Delta L_i}}{2 \, Y_{\gamma}} \, \text{Br}_X \, \text{Br}_L \langle \Gamma_S \rangle Y_S^{\text{eq}} \\ &\frac{H(m_X)}{z} \, \frac{dY_S}{dz} &= & -\langle \Gamma_S \rangle Y_S + \langle \Gamma_S \rangle Y_S^{\text{eq}} \left[ \text{Br}_L + \text{Br}_X \left( \frac{Y_X}{Y_X^{\text{eq}}} \right)^2 \right] \\ &\frac{H(m_X)}{z \, \eta} \, \frac{dY_{\Delta L_i}}{dz} &= & \frac{\epsilon_S}{2} \, \text{Br}_L \langle \Gamma_S \rangle \left[ Y_S + Y_S^{\text{eq}} \left( 1 - 2 \text{Br}_L - \text{Br}_X \left[ 1 + \frac{Y_X^2}{(Y_X^{\text{eq}})^2} \right] \right) \right] + 2s \, \epsilon_X \langle \sigma_{XX \leftrightarrow L_i \psi_i} \, v \rangle \left[ Y_X^2 - \frac{\xi \, Y_{\Delta L_i}}{Y_{\gamma}} \right] \\ &- \frac{\xi \, Y_{\Delta L_i}}{Y_{\gamma}} \left[ s \, \langle \sigma_{XX \leftrightarrow L_i \psi_i} \, v \rangle \langle Y_X^{\text{eq}} \rangle^2 + 2s [\langle \sigma_{L_i \psi_i \leftrightarrow L_i^\dagger \psi_i^\dagger} \, v \rangle + \langle \sigma_{L_i \psi_i \leftrightarrow L_j^\dagger \psi_j^\dagger}^{(i \neq j)} \, v \rangle Y_L^{\text{eq}} Y_\psi^{\text{eq}} \right. \\ &- \frac{2\xi \, Y_{\Delta L_i}}{Y_{\gamma}} s \, \langle \sigma_{L_i \psi_j \leftrightarrow L_j^\dagger \psi_i^\dagger} \, v \rangle Y_L^{\text{eq}} Y_\psi^{\text{eq}} \\ &- \frac{\xi \, Y_{\Delta L_i}}{Y_{\gamma}} \left[ s \, \langle \sigma_{X \psi_i \leftrightarrow X L_i^\dagger} \, v \rangle Y_X Y_\psi^{\text{eq}} + 2s \, \langle \sigma_{\psi_i \psi_i \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i^\dagger} \, v \rangle \langle Y_\psi^{\text{eq}} \rangle^2 + 2s \, \langle \sigma_{\psi_i \psi_j \leftrightarrow L_i^\dagger L_i$$

# Back-up slides: chemical potential relations

- $\bullet \quad \text{The } \psi \text{ mass: } \mu_{\psi} = -\mu_{\bar{\psi}}.$
- ② The SU(2) sphalerons:  $3\mu_Q + \mu_L = 0$ .
- **3** The up quark Yukawa:  $\mu_Q + \mu_H \mu_u = 0$ .
- **1** The down quark Yukawa:  $\mu_Q \mu_H \mu_d = 0$ .
- **1** The lepton Yukawa:  $\mu_L \mu_H \mu_E = 0$ .
- The  $\psi$  Yukawa:  $\mu_{\psi} \mu_H + \mu_{\chi} = 0$ .
- Hypercharge conservation:  $\mu_Q + 2\mu_u \mu_d \mu_L \mu_E + (\mu_\psi \mu_{\bar{\psi}}) \times (n_{\text{\tiny ab}}^{\text{eq}}/n_{\text{\tiny $\gamma$}}^{\text{eq}}) + 2\mu_H/3 = 0.$
- $\begin{array}{l} \textbf{ Onservation of generalized } B+\psi-L-\chi \text{ symmetry:} \\ 2\mu_Q+\mu_u+\mu_d-2\mu_L-\mu_E-\mu_\chi+2(\mu_\psi-\mu_{\bar\psi})\times(n_\psi^{\rm eq}/n_\gamma^{\rm eq})=0. \end{array}$

# Back-up slides: chemical potential solutions

$$\mu_{Q} = -\frac{1}{3}\mu_{L},$$

$$\mu_{u} = \frac{5 - 19r}{21 + 84r}\mu_{L},$$

$$\mu_{d} = -\frac{19 + 37r}{21 + 84r}\mu_{L},$$

$$\mu_{E} = \frac{3 + 25r}{7 + 28r}\mu_{L},$$

$$\mu_{H} = \frac{4 + 3r}{7 + 28r}\mu_{L},$$

$$\mu_{\chi} = -\frac{79 - 9r}{21 + 84r}\mu_{L},$$

$$\mu_{\psi} = \frac{13}{3 + 12r}\mu_{L},$$

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#### **Enhancement of asymmetry:**

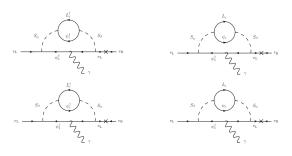
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#### Suppression of asymmetry:

- $\bullet$  Fraction of annihilations generating an asymmetry if  $1/\alpha$  , so  $Y_{\Delta B} \to Y_{\Delta B}/\alpha$
- Maximum allowed  $\epsilon$  is smaller because  $\lambda_L$  is smaller:  $\epsilon \to \epsilon/\sqrt{\alpha}$

### Constraints and signals: EDMs

- ullet Expect large CP phases to contribute to EDMs o CP problem
- New physics couples only to either LH or RH fields
  - Loops are helicity-preserving, so equal number of  $\lambda$  and  $\lambda^*$  insertions



$$\frac{d}{e} \sim \sum_{i} \operatorname{Im}(\lambda_{\alpha 1} \lambda_{\alpha i} \lambda_{\beta 1}^{*} \lambda_{\beta i}^{*} + \lambda_{\alpha 1} \lambda_{\alpha i}^{*} \lambda_{\beta 1} \lambda_{\beta i}^{*} + \text{c.c.}) = 0$$

- Vanishes when summed over permutations of internal lines!
  - ▶ No CP problem  $\rightarrow d/e < 10^{-30} \text{ e} \cdot \text{cm}$